

MATH327: StatMech and Thermo

Tuesday, 10 February 2026

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Something to consider

A micro-canonical system in thermodynamic equilibrium
adopts all possible micro-states with equal probability.

How can this explain the smooth and stable distribution
of the $\sim 10^{25}$ molecules that compose the air in this room?

Recap

Statistical ensembles

↳ Stochastically sampling micro-states w_i with prob P_i

Conserve $E, N \rightarrow$ micro-~~st~~ canonical ensemble
"first law"

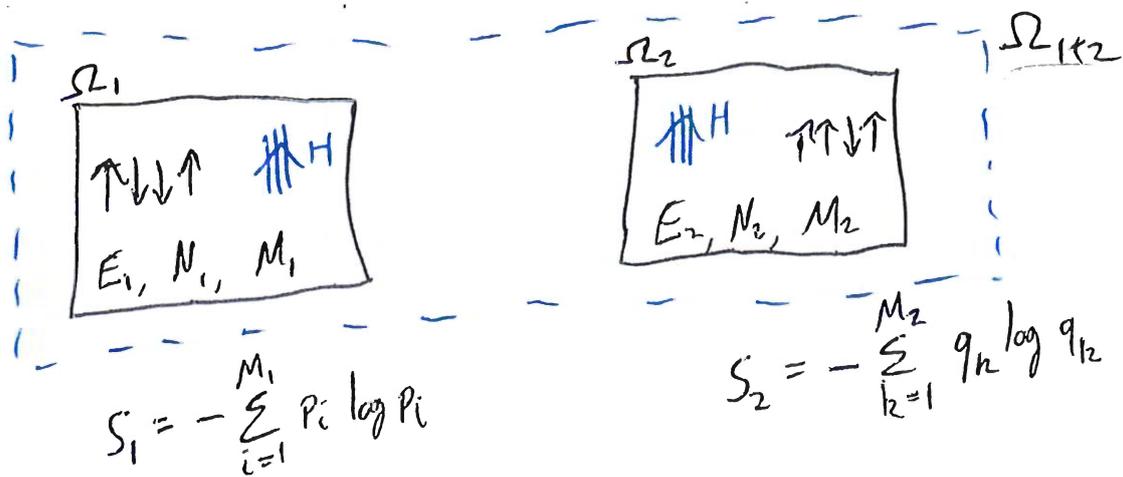
Entropy $S = - \sum_{i=1}^M P_i \log P_i$

Thermodynamic equilibrium

↳ Micro-canonical all P_i equal $\rightarrow S = \log M$

Today

Second law \rightarrow generalize therm equil. based on entropy



What is S_{1+2} ?

For each $w_i \in \Omega_1$, M_2 micro-states from Ω_2 .

$\therefore M_{1+2} = M_1 M_2$ micro-states with prob. $p_i q_k$

Sanity check: $\sum_{M_{1+2}} p_i q_k = \sum_{i=1}^{M_1} \sum_{k=1}^{M_2} p_i q_k = \left(\sum_i p_i \right) \left(\sum_k q_k \right) = 1 \checkmark$

$$S_{1+2} = - \sum_{i,k} p_i q_k \log(p_i q_k) = \sum_{i,k} p_i q_k (\log p_i + \log q_k)$$

$$= - \sum_i p_i \log p_i \left(\sum_k q_k \right) - \left(\sum_i p_i \right) \sum_k q_k \log q_k = S_1 + S_2$$

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Extensive quantity adds up across independent subsystems

$$S_{1+2} = S_1 + S_2$$

$$N_{1+2} = N_1 + N_2$$

$$E_{1+2} = E_1 + E_2$$

Intensive quantity independent of extent of system
 temperature, pressure, density

$$M_{1+2} = M_1 M_2 \quad \text{neither intensive nor extensive}$$

IF each of Ω_1, Ω_2 in therm. equil.

$$P_i = \frac{1}{M_1}$$

$$q_k = \frac{1}{M_2}$$

$$S_1 = \log M_1$$

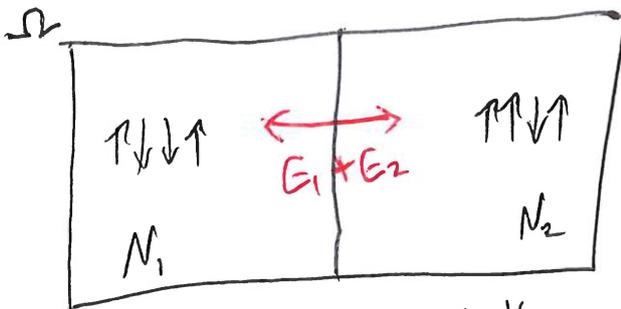
$$S_2 = \log M_2$$

$$P_i q_k = \frac{1}{M_1} \cdot \frac{1}{M_2} = \frac{1}{M_{1+2}} \rightarrow \text{also in equil.}$$

$$S_{1+2} = \log M_{1+2} = \log M_1 + \log M_2 = S_1 + S_2 \quad \checkmark$$

Now put subsystems in thermal contact

↳ exchange energy, not particle



wait for equilibrium

Total $E = E_1 + E_2$ still conserved $\rightarrow \Omega$ is micro-canonical
 subsystems are not

How many micro-states M for overall Ω ?

Use energy conservation: $\left. \begin{array}{l} e_1 \text{ from } N_1 \\ E - e_1 \text{ from } N_2 \end{array} \right\} M_{e_1} = M_{e_1}^{(1)} M_{E-e_1}^{(2)}$

$$\text{Overall } M = \sum_{e_1} M_{e_1}^{(1)} M_{E-e_1}^{(2)}$$

$e_1 = E_1$ accounts for all $M_{1+2} = M_1 M_2 = M_{E_1}^{(1)} M_{E_2}^{(2)}$

$$M = M_{1+2} + \sum_{e_1 \neq E_1} M_{e_1}^{(1)} M_{E-e_1}^{(2)} \geq M_{1+2}$$

→ $M > M_{1+2}$ unless $\{E_1, E_2\}$ only possible distribution
very special case

Result $S = \log M \geq \log (M_{1+2}) = S_{1+2} = S_1 + S_2$

Second law of thermodynamics

When isolated subsystems in therm. equil.
are brought into thermal contact
entropy $S \geq S_1 + S_2$ can never decrease

Increases except in very special cases

More generally, entropy never decreases as time passes
(generally increases)

→ System (with finite M) is in
thermodynamic equilibrium
if its entropy is maximal

Holds for any stat. ensemble

Let's check micro-canonical

Maximize $S = - \sum_i p_i \log p_i$ with conserved E, N
and $\sum_i p_i = 1$

Use Lagrange multiplier

$$\bar{S}(\lambda) = S + \lambda (\sum_i p_i - 1) = - \sum_i p_i \log p_i + \lambda (\sum_i p_i - 1)$$

1) Maximize $\bar{S}(\lambda)$ w.r.t. p_k

2) Impose $\sum_i p_i = 1 \rightarrow \frac{\partial \bar{S}}{\partial \lambda} = 0$ so $\max \bar{S} \leftrightarrow \max S$
 $\bar{S}(\lambda=0)$

$$\frac{\partial \bar{S}}{\partial P_h} = 0 = \frac{\partial}{\partial P_h} \left[- \sum_i P_i \log P_i + \lambda \left(\sum_i P_i - 1 \right) \right]$$

$$= - \log P_h - \frac{P_h}{P_h} + \lambda$$

$$\log P_h = \lambda - 1 \quad \rightarrow \quad P_h = \exp(\lambda - 1) = \text{const.} = \frac{1}{M} \quad \checkmark$$

$$S = \log M$$

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Derived micro-canonical def. of therm. equil.
from entropy maximization