

# MATH327: Statistical Physics

Tuesday, 27 February 2024

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## Something to consider

Systems governed by the canonical ensemble

can have different internal energy  $E_i$  for each micro-state  $\omega_i$ .

In the very special case that all micro-states have the same energy  $E$ , will we observe behaviour similar to the micro-canonical ensemble?

Recap

Micro-canonical temperature and heat exchange

Canonical ensemble - fix  $T$  via thermal reservoir

$$\Omega_{\text{tot}} = \Omega \otimes \Omega_{\text{res}} \quad \text{micro-canonical}$$

Replica trick: Suppose  $\Omega_{\text{res}}$  is  $R-1$  copies of  $\Omega$   
 $\Rightarrow$

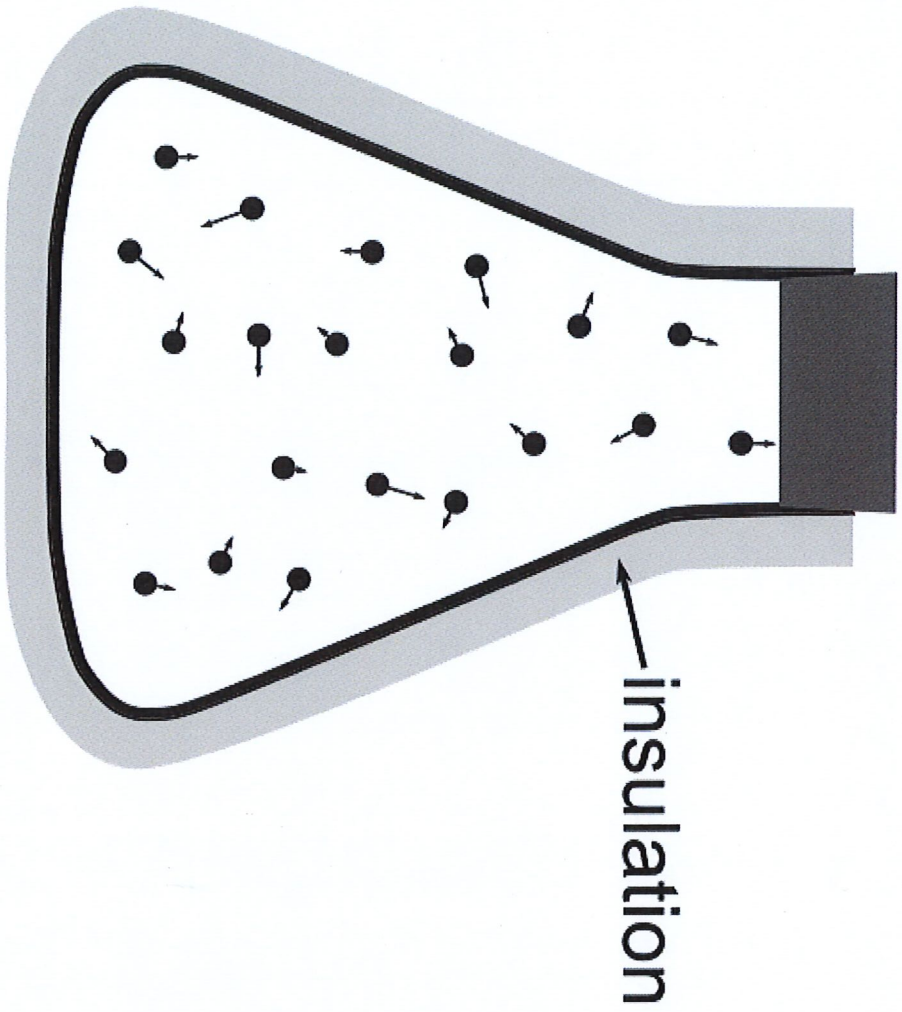
Number of micro-states  $M_{\text{tot}} = \frac{R!}{n_1! n_2! n_3! \dots n_M!}$

Entropy in therm. equil.  $S_{\text{tot}} = \log M_{\text{tot}} = \log(R!) - \sum_{i=1}^M \log(n_i!)$

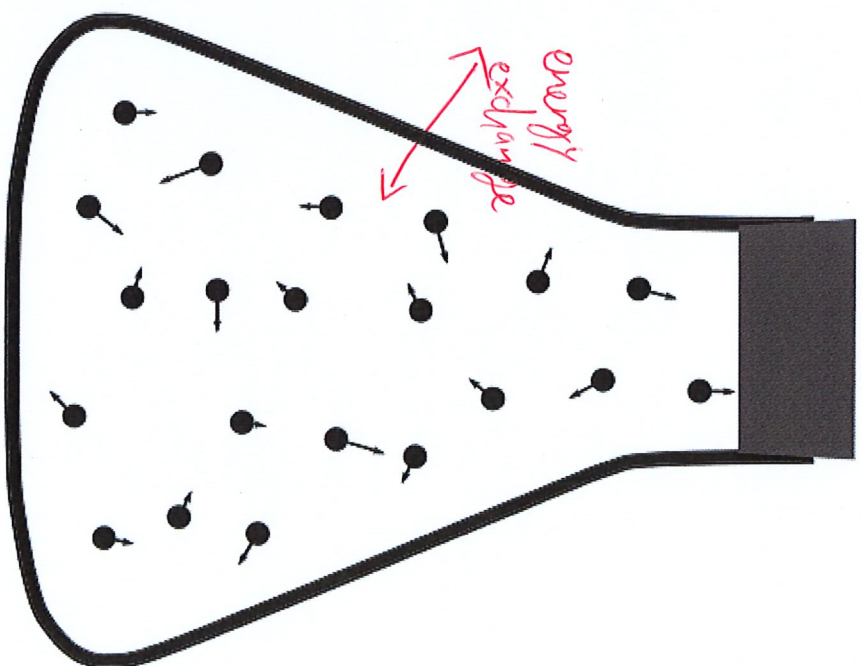
Assume all  $n_i \gg 1$  and approx.  $\log(n!) \approx n \log n - n$

$$S_{\text{tot}} \approx R \log R - R - \sum_i (n_i \log n_i - n_i) \quad n_i = p_i R$$

$$= R \log R - R \sum_i p_i (\log p_i + \log R) = -R \sum_i p_i \log p_i$$



**Microcanonical**  
(const.  $N$   $E$ )



**Canonical**  
(const.  $N$   $T$ )

Maximize entropy with constraints  $\sum_i p_i = 1$

$$R \sum_i p_i E_i = E_{\text{tot}}$$

$$\bar{S} = -R \sum_i p_i \log p_i + \alpha \left( \sum_i p_i - 1 \right) - \beta \left( R \sum_i p_i E_i - E_{\text{tot}} \right)$$

$$\frac{\partial \bar{S}}{\partial p_k} = 0 = -R (\log p_k + 1) + \alpha - \beta R E_k$$

$$\log p_k = -\frac{1}{R} + \frac{\alpha}{R} - \beta E_k$$

$$p_k = \exp \left[ -\left(1 - \frac{\alpha}{R}\right) - \beta E_k \right] = \frac{\exp(-\beta E_k)}{\exp\left(1 - \frac{\alpha}{R}\right)} = \frac{1}{Z} e^{-\beta E_k}$$

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Impose constraints

$$\sum_i p_i = 1 = \frac{1}{Z} \sum_i e^{-\beta E_i} \rightarrow Z(\beta) = \sum_i e^{-\beta E_i}$$

partition function  
(fundamental)

$$R \sum_i p_i E_i = E_{\text{tot}}$$

$$S_{\text{tot}} = -R \sum_i p_i \log p_i = -R \sum_i p_i \log \left( \frac{1}{Z} e^{-\beta E_i} \right)$$

$$= R \log Z + R \beta \sum_i p_i E_i$$

$$= R \log Z + \beta E_{\text{tot}}$$

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$$\frac{1}{T} = \frac{\partial S}{\partial E} = \beta + E \frac{\partial \beta}{\partial E} + R \left( \frac{1}{Z} \frac{\partial Z}{\partial \beta} \frac{\partial \beta}{\partial E} \right)$$

$$\frac{1}{Z} \frac{\partial Z}{\partial \beta} = \frac{1}{Z} \sum_i \frac{\partial}{\partial \beta} e^{-\beta E_i} = -\frac{1}{Z} \sum_i E_i e^{-\beta E_i}$$

$$= -\sum_i p_i E_i = -\frac{E_{\text{tot}}}{R}$$

$$\frac{1}{T} = \beta + E \frac{\partial \beta}{\partial E} + R \left( -\frac{E}{R} \right) \frac{\partial \beta}{\partial E} = \beta$$

We have derived the Gibbs distribution

$$P_i = \frac{1}{Z} e^{-E_i/T} \quad Z = \sum_i e^{-E_i/T} \quad \frac{1}{T} = \beta$$

Boltzmann factor

$P_i$  are probability system  $\Omega$  adopts micro-state  $w_i$   
(Equal  $E_i \rightarrow$  equal  $P_i$ ) with energy  $E_i$

Reservoir unknowable apart from fixing  $T$

Derived quantities

Expectation value of internal energy

$$\langle E \rangle(T) = \sum_i P_i E_i = \frac{1}{Z} \sum_i E_i e^{-\beta E_i} = -\frac{\partial}{\partial \beta} \log Z$$

How does  $\langle E \rangle$  depend on  $T$ ?

$$c_v = \frac{\partial \langle E \rangle}{\partial T} \geq 0$$

Heat capacity

(HW: fluctuation-response)

Higher temperature  
 $\rightarrow$  larger internal energy

Entropy

$$S = -\sum_i P_i \log P_i = -\sum_i P_i \log \left( \frac{1}{Z} e^{-\beta E_i} \right)$$

$$= \log Z + \beta \sum_i P_i E_i = \log Z + \frac{\langle E \rangle}{T}$$

$$\langle E \rangle = T \cdot S - T \log Z = T \cdot S + F$$

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Helmholtz Free energy

$$F(T) = -T \log Z$$

$$Z = e^{-F/T}$$

$$P_i = \frac{1}{Z} e^{-E_i/T} = e^{(F-E_i)/T}$$