

Thu 20 Apr

669067

Plan

Photon polarization

Dark matter in CMB

Einstein solid & Debye solid

Heat capacity $c_v = \frac{\partial E}{\partial T} \Big|_V$

Ideal gas $\langle E \rangle = \frac{3}{2} NT \rightarrow c_v = \frac{3}{2} N \quad \times$

Di^fferentiable spins $c_v = \frac{N\beta^2 H^2}{\cosh^2(\beta H)} > 0$

Low-T $\beta \rightarrow \infty \quad \cosh^2(\beta H) \approx e^{2\beta H} \rightarrow \infty \quad c_v \rightarrow 0 \quad \checkmark$
(third law)

High-T $\beta \rightarrow 0 \quad \cosh^2(\beta H) \rightarrow 1 \quad c_v \rightarrow 0 \quad \times$

Einstein solid

Fix atoms in solid by interacting w/ neighbours
↳ energy in oscillators rather than atoms

Quantized energies for oscillators $\epsilon_i = 0, \hbar\omega, 2\hbar\omega, \dots$

Micro-canonical: Fix N oscillators sharing K units of energy

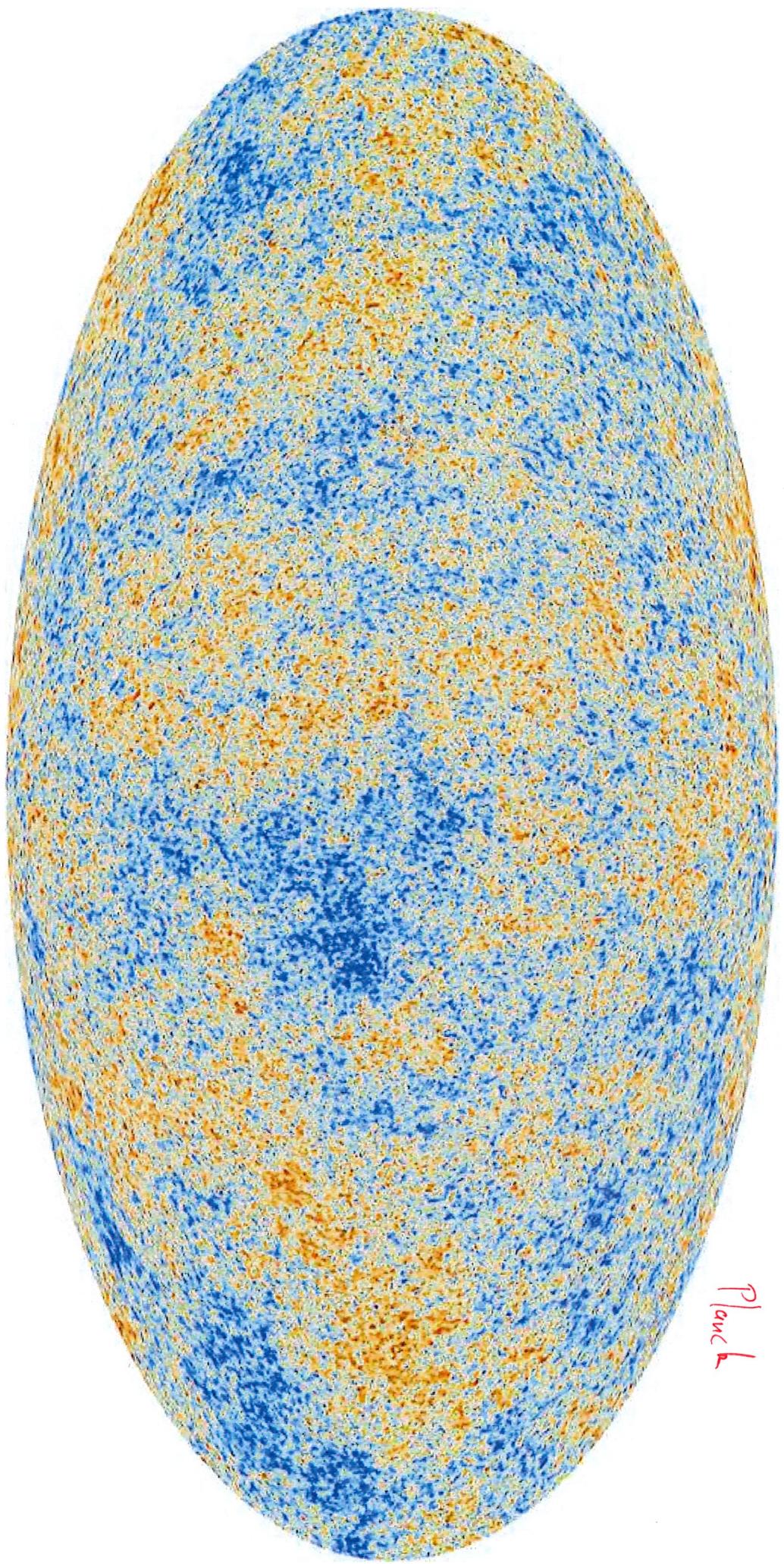
$$1) \frac{1}{T} = \frac{\partial S}{\partial E} = \frac{2}{\partial E} \log M$$

$$K = \sum_{i=1}^N k_i$$

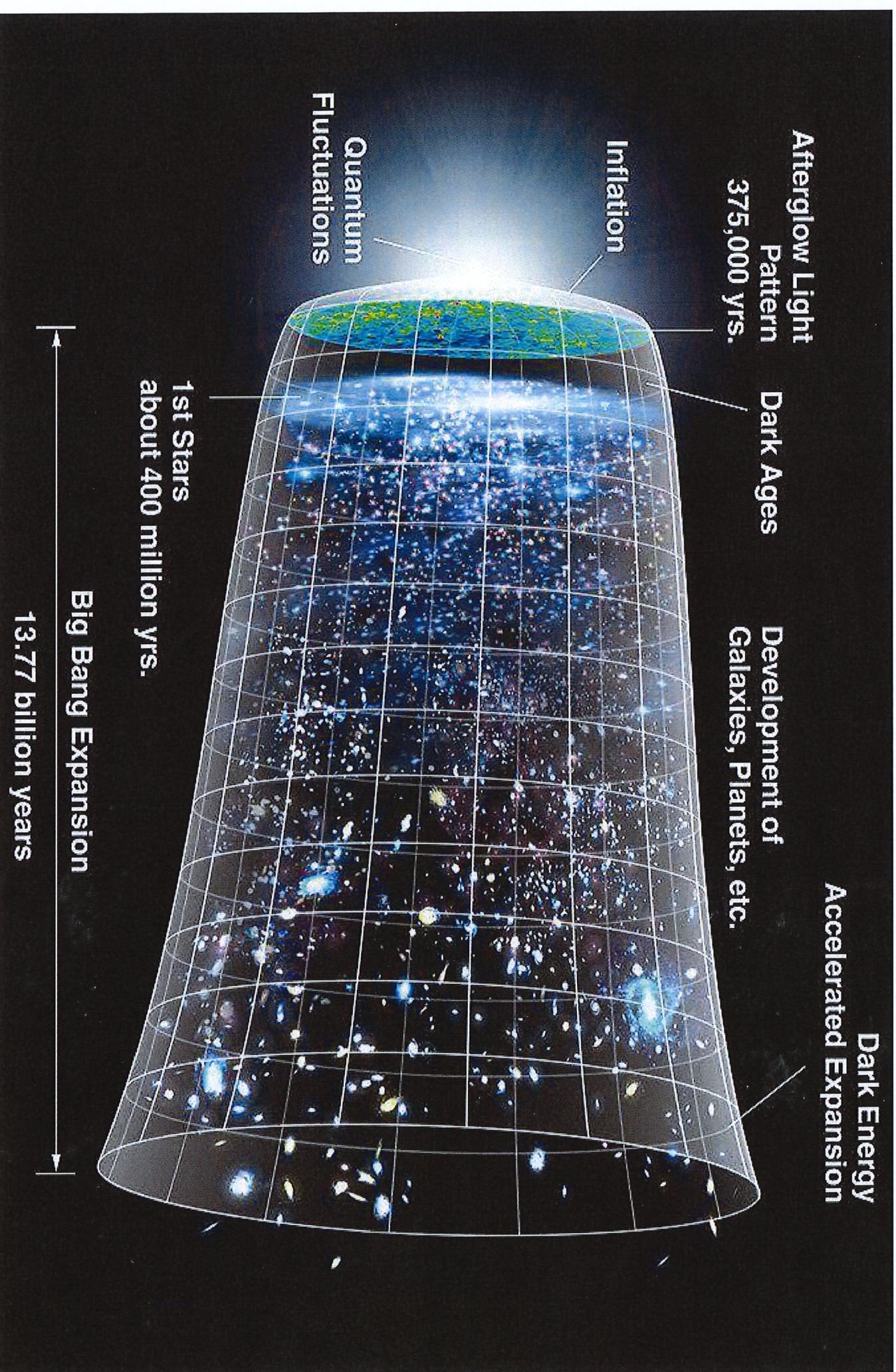
$$2) E(T)$$

$$E = \sum_i \hbar \epsilon_i \text{th} \omega$$

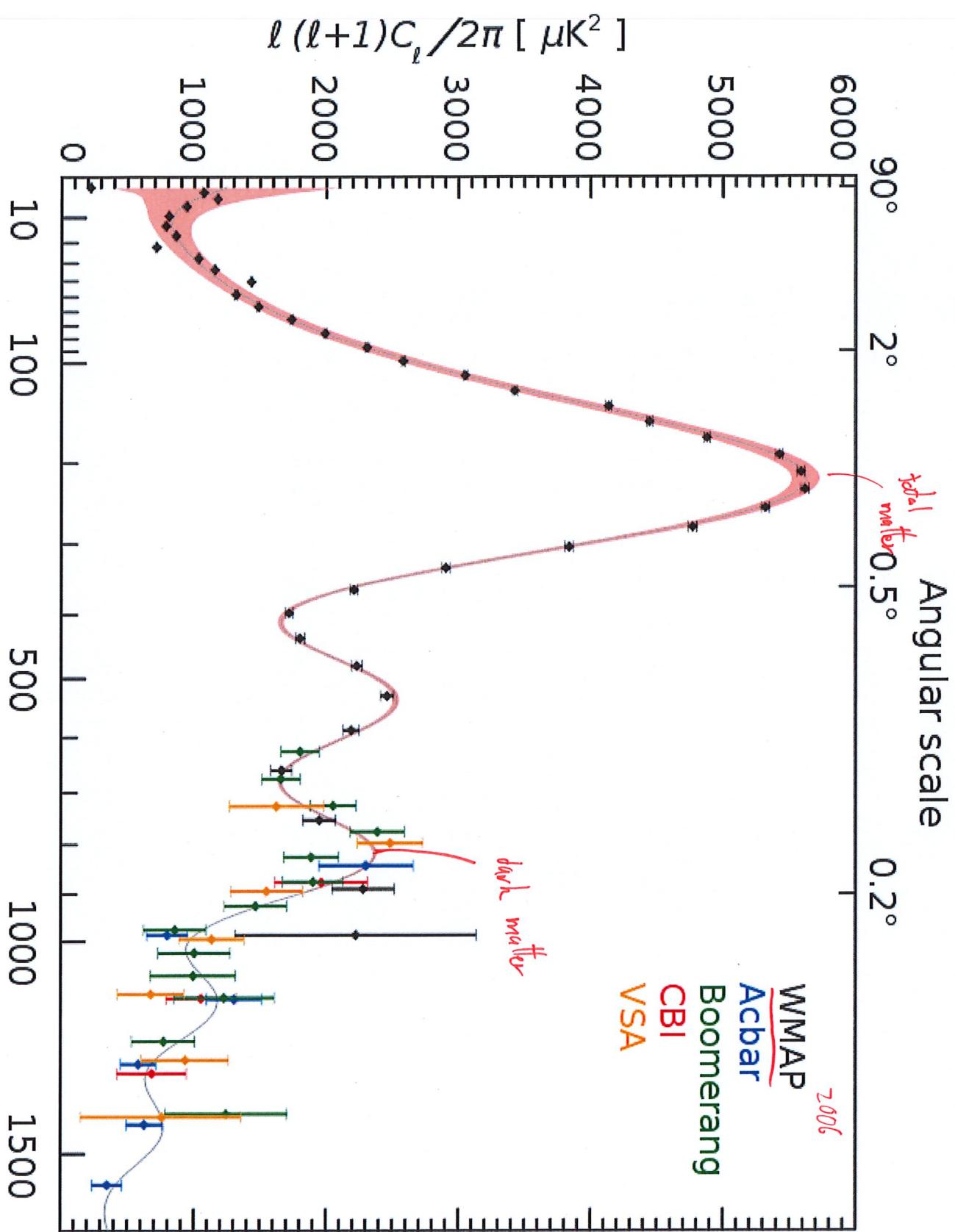
$$3) c_v = \frac{\partial E}{\partial T}$$



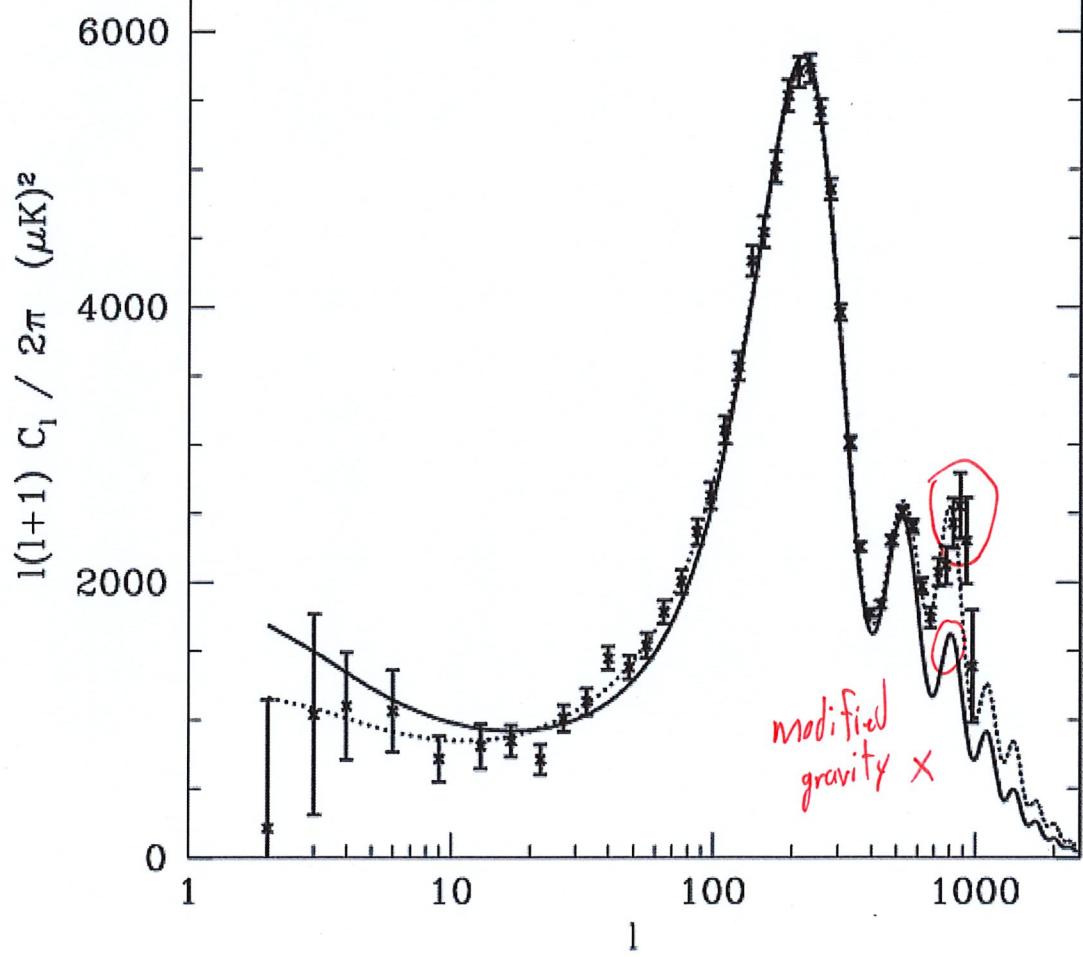
Planck

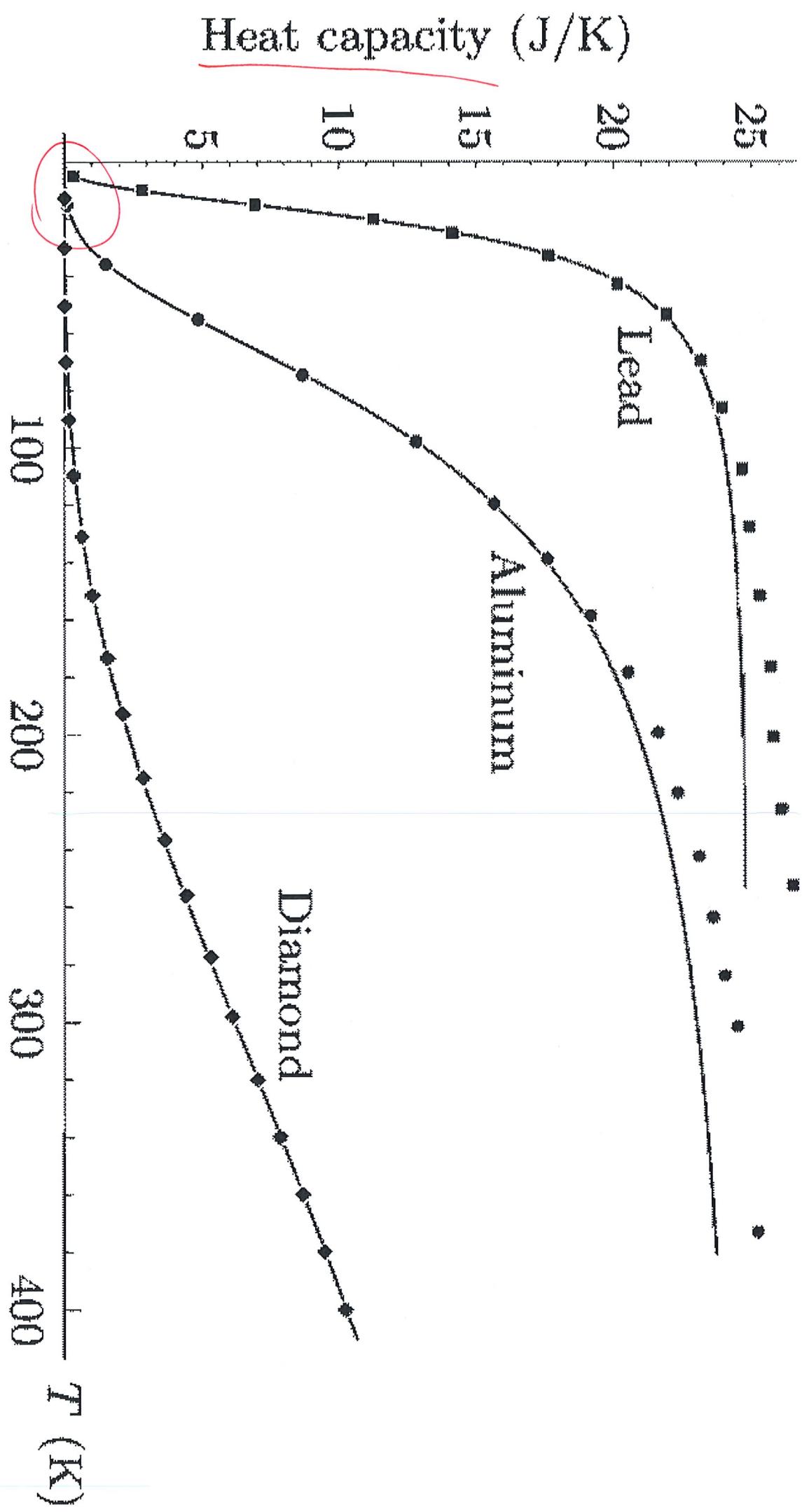


commons.wikimedia.org/wiki/File:CMB_Timeline_300-nl_WMAP.jpg

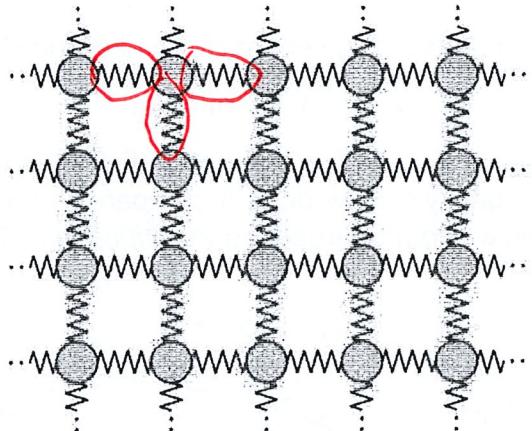


arXiv:0903.3602





their layout. In this two-dimensional square lattice, N oscillators would correspond to $N/2$ atoms in the solid. In a three-dimensional cubic lattice, N oscillators would correspond to $N/3$ atoms.



Our goal is to compute the heat capacity for an Einstein solid. Let's begin by working in terms of the micro-canonical ensemble, fixing the total energy

$$E = \sum_i \varepsilon_i = \sum_i k_i \hbar \omega \equiv K \hbar \omega$$

where $K \equiv \sum_i k_i$ is the integer number of energy ‘units’ available to be distributed among the N oscillators. Each different way of distributing these K units of energy among the N (distinguishable) oscillators defines a unique micro-state.

- What is the total number of micro-states in terms of N and K ? Check your result for a minimal three-oscillator system when it has $K = 0, 1, 2$ or 3 units of energy.
- Now consider $K \gg 1$ and $N \gg 1$, so that we can apply Stirling’s formula and also approximate $N-1 \approx N$ and $K-1 \approx K$. What is the corresponding entropy?
- What is the temperature of the Einstein solid? Is it a ‘natural’ system with $T > 0$?
- To find the heat capacity, we need to find the energy in terms of the temperature, then differentiate. What is the resulting heat capacity of the Einstein solid, in terms of $x \equiv \hbar \omega / T$? How does it compare to the experimental data points we began by considering?



While the Einstein solid provides a significantly better description of the experimental data, room for improvement is still possible...



$N=3$

$$K=0 \quad M=1$$

$$K=1 \quad M=3$$

$$K=2 \quad M=3+3=6$$

$$K=3 \quad M=3+6+1=10$$

}

$$M = \frac{(K+N-1)!}{K! (N-1)!} = \binom{K+N-1}{K}$$