

Mon 17 Apr

497 472

Logistics

HW2 due 3 May

Exam 14:30 - 16:30 Friday 26 May

Big-picture review

Prob. spaces $\rightarrow$ stat. ensembles
micro-canonical
canonical (applications to ideal gases
therm. cycles)
grand-canonical

Plan

Grand-canonical applications to quantum gases

Quantum statistics: Micro-states defined via
energy level occupation numbers

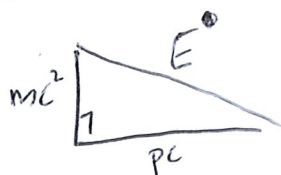
Bosons $n_l = 0, 1, 2, \dots \rightarrow$ Bose-Einstein
$$Z_g^{(BE)} = \prod_{l=0}^{\infty} \frac{1}{1 - e^{-\beta(E_l + \mu)}}$$
 $\mu < E_l$

Fermions $n_l = 0, 1 \rightarrow$ Fermi-Dirac
$$Z_g^{(FD)} = \prod_{l=0}^{\infty} (1 + e^{-\beta(E_l + \mu)})$$

Classical limit: high temperatures w/ large negative chem. pot.
 $-\mu \gg T \gg E_l$

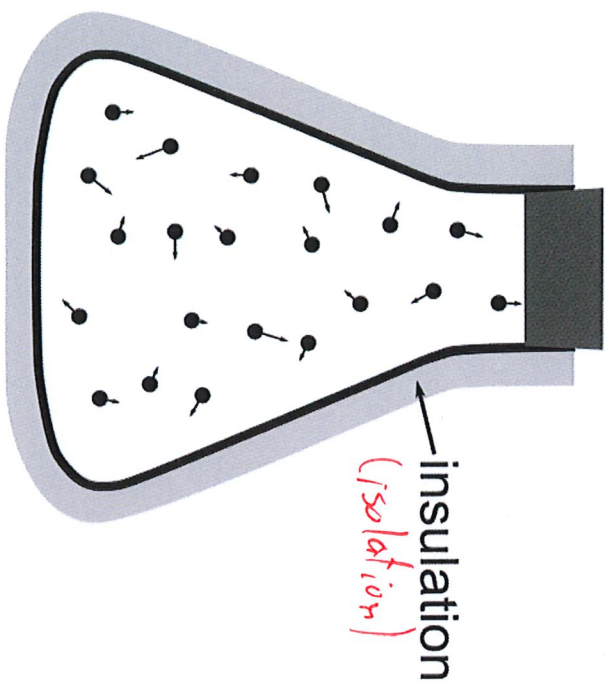
Need energy levels and E_l

In general $E^2 = (mc^2)^2 + (pc)^2$ speed of light (unit conversion)

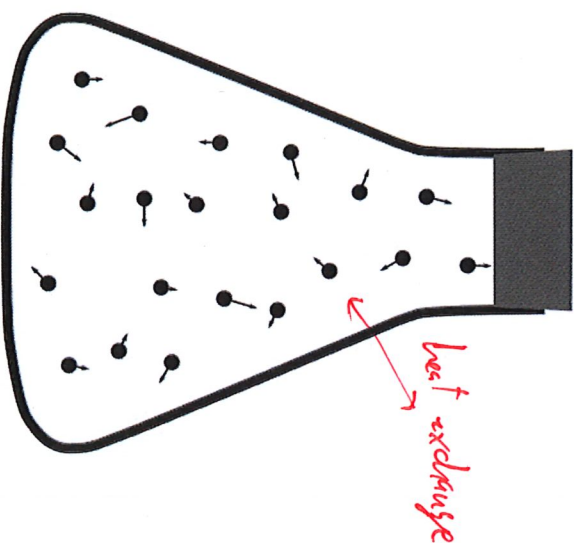


mass energy & kinetic energy
(Fixed)

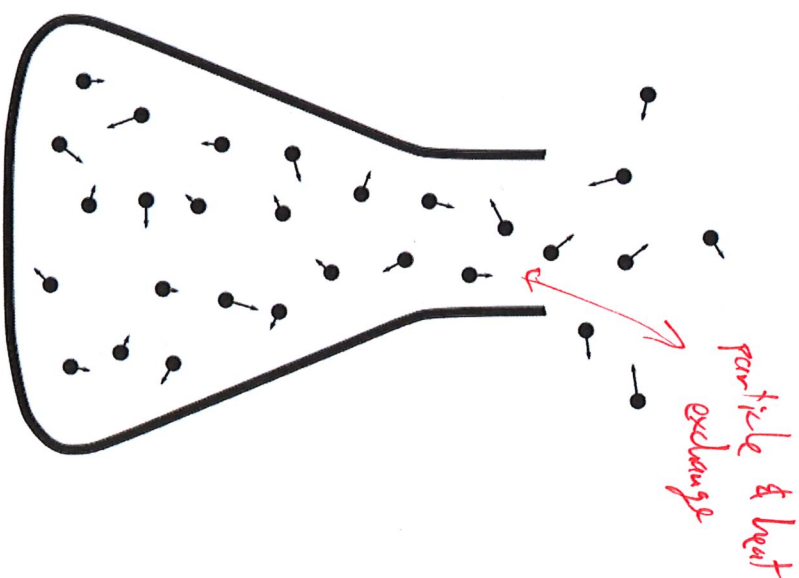
$$p^2 = p_x^2 + p_y^2 + p_z^2$$



Microcanonical
(const. N E)



Canonical
(const. N T)



Grand Canonical
(const. μ T)
chemical potential

$m=0$ or $p \gg mc \rightarrow$ ultra-relativistic $E=pc$

Non-relativistic limit $p \ll mc$

$$E = mc^2 \sqrt{1 + \frac{(pc)^2}{(mc^2)^2}} = mc^2 \left(1 + \frac{p^2}{2m^2c^2} + \mathcal{O}\left(\frac{p^4}{m^4c^4}\right) \right)$$
$$= \underbrace{mc^2}_{\text{irrelevant}} + \frac{p^2}{2m} + \mathcal{O}\left(\frac{p^4}{m^3c^2}\right)$$

page 103

$$E = \frac{p^2}{2m} = \frac{\hbar^2 \pi^2}{2mL^2} (k_x^2 + k_y^2 + k_z^2) \quad \text{in } L^3 \text{ volume}$$

now with $k_{x,y,z} = 1, 2, 3, \dots$ (Heisenberg)

$$\begin{array}{ll} \text{Quantum ground state} & \vec{k} = (1,1,1) \rightarrow E = 3\varepsilon \\ \text{excited state} & \vec{k} = (2,1,1) \rightarrow E = 6\varepsilon \\ & \vec{k} = (2,2,1) \rightarrow E = 9\varepsilon \\ & \vec{k} = (3,1,1) \rightarrow E = 11\varepsilon \\ & \vec{k} = (2,2,2) \rightarrow E = 12\varepsilon \end{array} \left. \vphantom{\begin{array}{l} \vec{k} = (2,1,1) \\ \vec{k} = (2,2,1) \\ \vec{k} = (3,1,1) \\ \vec{k} = (2,2,2) \end{array}} \right\} \text{each } 3\times \text{ degenerate}$$
$$\varepsilon = \frac{\hbar^2 \pi^2}{2mL^2}$$

page 103

Consider gas of photons (bosonic) with $m=0$ $E_{ph} = pc = \hbar c \frac{\pi}{L} k$

Photons are quanta of electromagnetic waves (light)

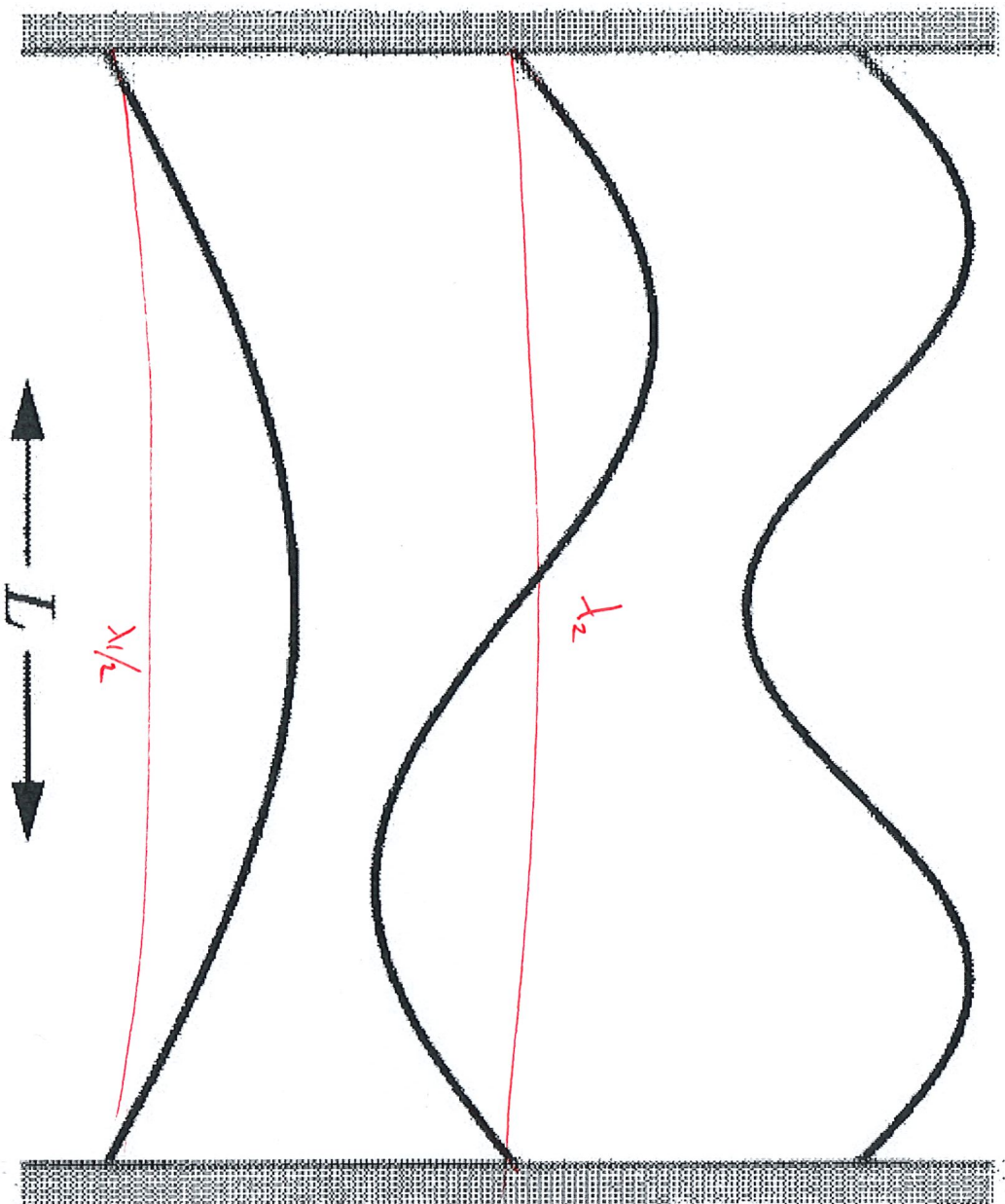
Speed of light $c = \frac{\lambda \omega}{2\pi}$ from wavelength λ
angular frequency $\omega = 2\pi f$

Volume $V=L^3$ quantizes frequency like momenta

$$L = k_x \left(\frac{\lambda}{2} \right) \rightarrow \lambda = \frac{2L}{k_x} \quad k_x = 1, 2, \dots$$

$$\omega = \frac{2\pi c}{\lambda} = c \frac{\pi}{L} k_x = \left(\frac{c}{\hbar} \right) p_x$$

$$E_{ph} = \hbar \omega$$

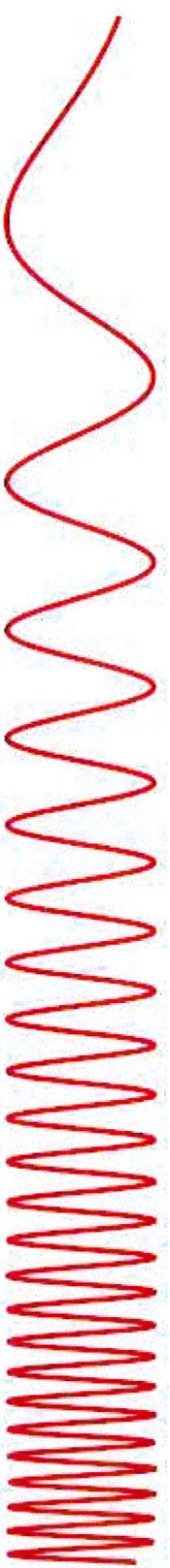


$$\lambda_1 = 2L$$

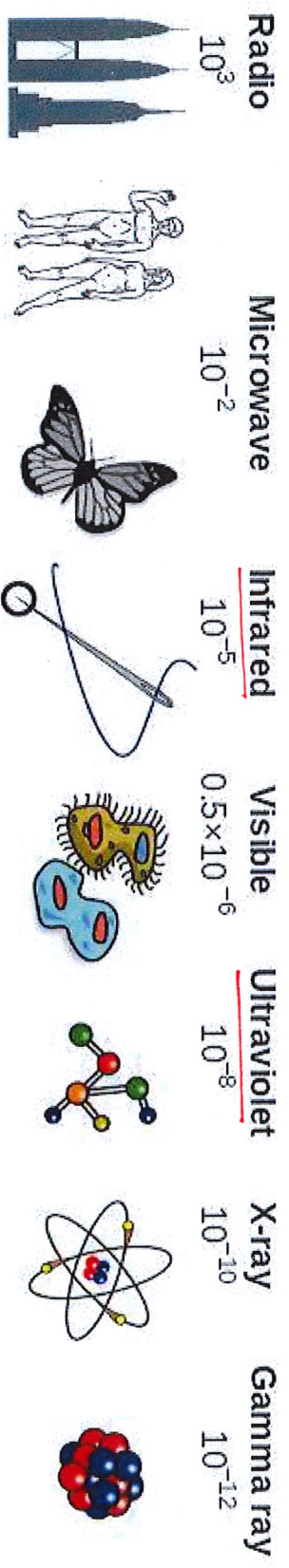
$$\lambda_2 = \frac{2L}{2}$$

$$\lambda_3 = \frac{2L}{3}$$

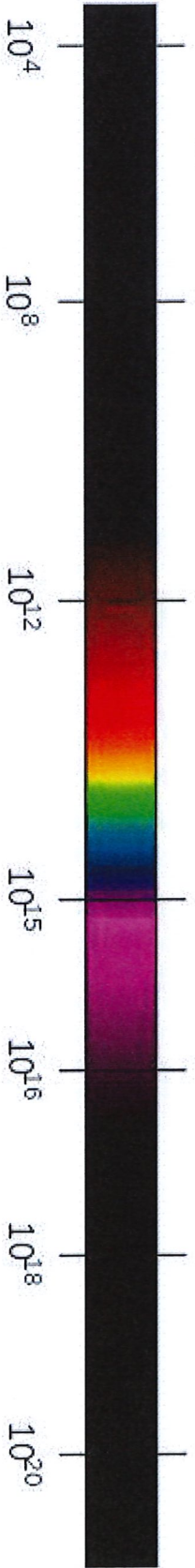
Penetrates Earth's Atmosphere?



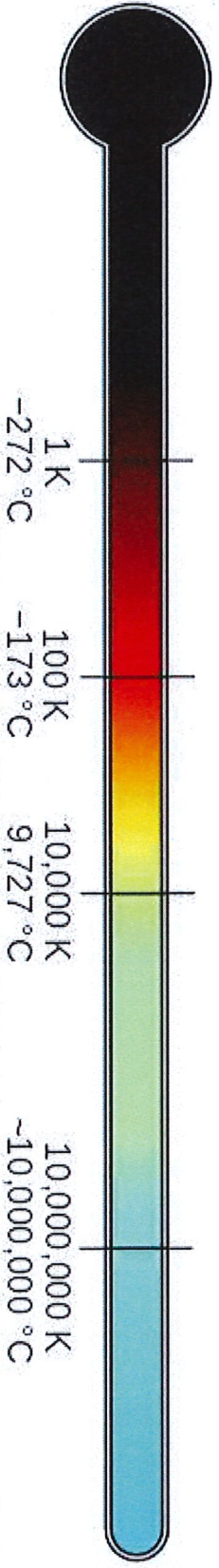
Radiation Type
Wavelength (m)
Approximate Scale of Wavelength



Frequency (Hz)



Temperature of objects at which this radiation is the most intense wavelength emitted



commons.wikimedia.org/wiki/File:EM_Spectrum_Properties_edit.svg

Photon gas grand-canonical potential

$$\Phi_{ph} = T \sum_{\vec{k}} \log(1 - e^{-\beta(E_{ph} - \mu)}) = 2T \sum_{\vec{k}} \log(1 - e^{-\beta(\hbar\omega - \mu)})$$

two polarizations per $\vec{k}$

Simplification: Photons easy to create & absorb

Little energy to photon #

$$\mu = \left. \frac{\partial E}{\partial N} \right|_S \approx 0$$

$$E_{ph} = \hbar\omega > 0 \rightarrow \mu < E_{ph} \checkmark$$

Simplification: Integrate over closely spaced momenta

$$\Phi_{ph} \approx 2T \int \log(1 - e^{-\beta\hbar\omega}) d\hat{k}_x d\hat{k}_y d\hat{k}_z$$

$\omega \propto k \rightarrow$ spherical coordinates
positive $\rightarrow$ single octant

$$\int_0^\infty d\hat{k}_x \int_0^\infty d\hat{k}_y \int_0^\infty d\hat{k}_z = \int_0^\infty \hat{k}^2 d\hat{k} \int_0^{\pi/2} \sin\theta d\theta \int_0^{\pi/2} d\phi = \frac{\pi}{2} \int_0^\infty \hat{k}^2 d\hat{k}$$

plug in $k = \omega \left(\frac{L}{c\pi} \right)$

$$\Phi_{ph} \approx \pi T \int_0^\infty \hat{k}^2 \log(1 - e^{-\beta\hbar\omega}) d\hat{k}$$

$$= \frac{VT}{c^3 \pi^2} \int_0^\infty \omega^2 \log(1 - e^{-\beta\hbar\omega}) d\omega$$

Internal energy

$$\langle E \rangle = -T^2 \frac{\partial}{\partial T} \left(\frac{\Phi}{T} \right) + \mu \langle N \rangle = \frac{\partial}{\partial \beta} (\beta \Phi)$$

$$\langle E \rangle_{ph} = \frac{V}{c^3 \pi^2} \int_0^\infty \omega^2 \frac{\partial}{\partial \beta} \log(1 - e^{-\beta\hbar\omega}) d\omega$$

$$= \frac{V}{c^3 \pi^2} \int_0^\infty \frac{\omega^2 (-e^{-\beta\hbar\omega}) (-\hbar\omega)}{1 - e^{-\beta\hbar\omega}} d\omega = \frac{V\hbar}{c^3 \pi^2} \int_0^\infty \frac{\omega^3}{e^{\beta\hbar\omega} - 1} d\omega$$

page 106

Energy density $\frac{\langle E \rangle_{ph}}{V} = \int_0^\infty P(\omega) d\omega$
spectral density

For photons $P(\omega) = \left(\frac{h}{c^3 \pi^2} \right) \frac{\omega^3}{e^{\beta h \omega} - 1}$ is Planck spectrum

Equivalently write as $P(\lambda)$ with $\lambda = \frac{2\pi c}{\omega}$ $\omega = \frac{2\pi c}{\lambda}$ $d\omega = -\frac{2\pi c}{\lambda^2} d\lambda$

$$\frac{\langle E \rangle_{ph}}{V} = \frac{h}{c^3 \pi^2} \int_0^\infty \frac{(2\pi c/\lambda)^3}{e^{\frac{2\pi \beta h c}{\lambda}} - 1} \left(-\frac{2\pi c}{\lambda^2} \right) d\lambda = 16\pi^2 h c \int_0^\infty \frac{d\lambda}{\lambda^5 (e^{\frac{2\pi \beta h c}{\lambda}} - 1)}$$

page 106

$$P(\lambda) = \left(\frac{16\pi^2 h c}{\lambda^5} \right) \frac{1}{e^{\frac{2\pi \beta h c}{\lambda}} - 1}$$

UV: $\lambda \rightarrow 0$

Exponential factor dominates $\frac{1}{\lambda^5}$ term

IR: λ large compared to $h c$

$$e^{\frac{2\pi \beta h c}{\lambda}} - 1 \approx \frac{2\pi \beta h c}{\lambda} \rightarrow P(\lambda) = \left(\frac{16\pi^2 h c}{\lambda^5} \right) \left(\frac{\lambda}{2\pi \beta h c} \right)$$

$$= \frac{8\pi T}{\lambda^4}$$

also for small β
(high temperatures)

classical
Rayleigh-Jeans spectrum

Classical spectrum diverges $\sim \frac{1}{\lambda^4}$ as $\lambda \rightarrow 0$

"ultraviolet catastrophe" $\rightarrow$ quantum

Max of Planck $P(\lambda)$ at shorter λ for higher temperatures

Peak in λ of visible light for $T \approx 5000$ K

because sunlight has $T \approx 5778$ K

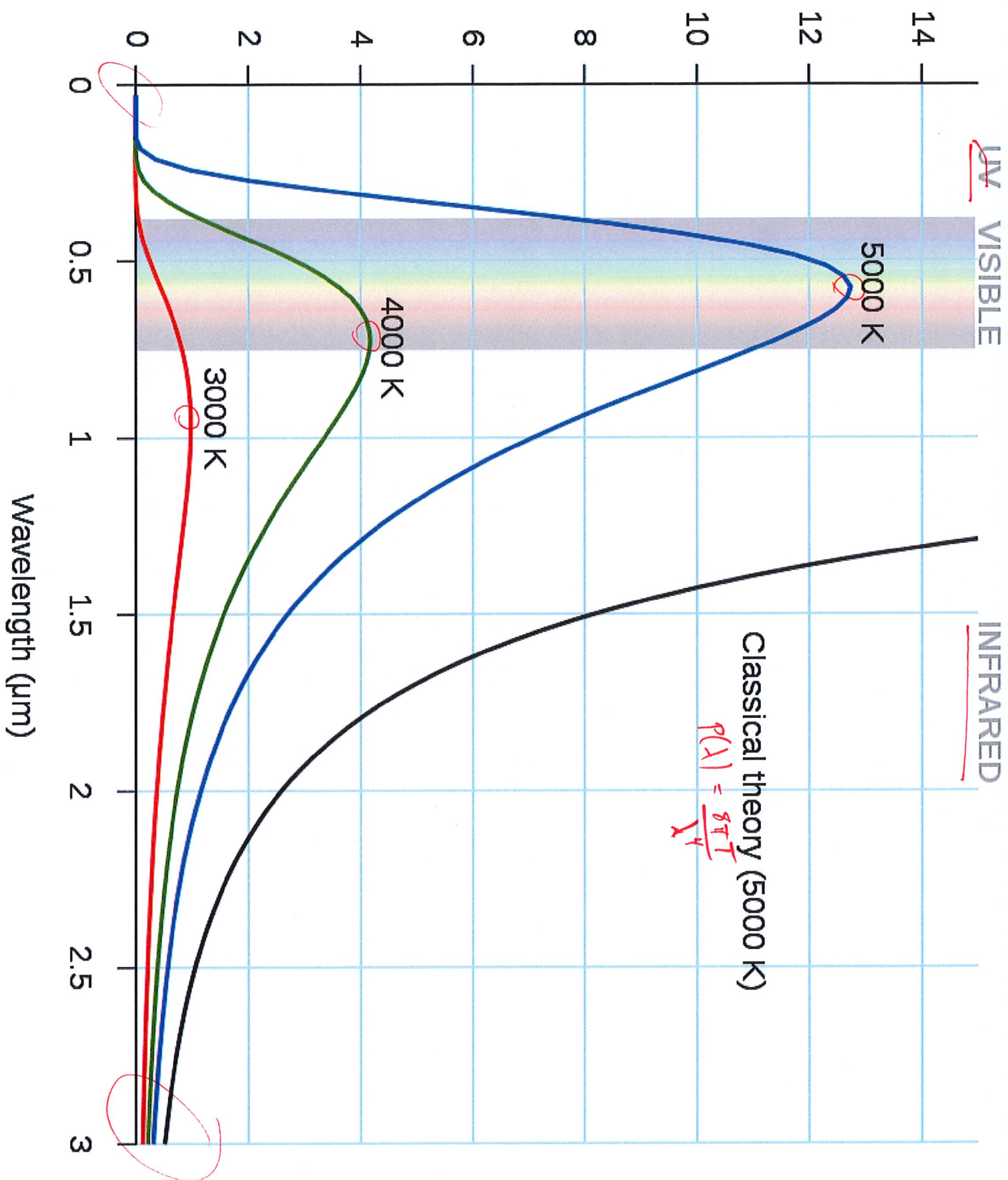
Determine effective surface temperature by fitting $P(\lambda)$

Red stars $\rightarrow T \approx 3500$ K

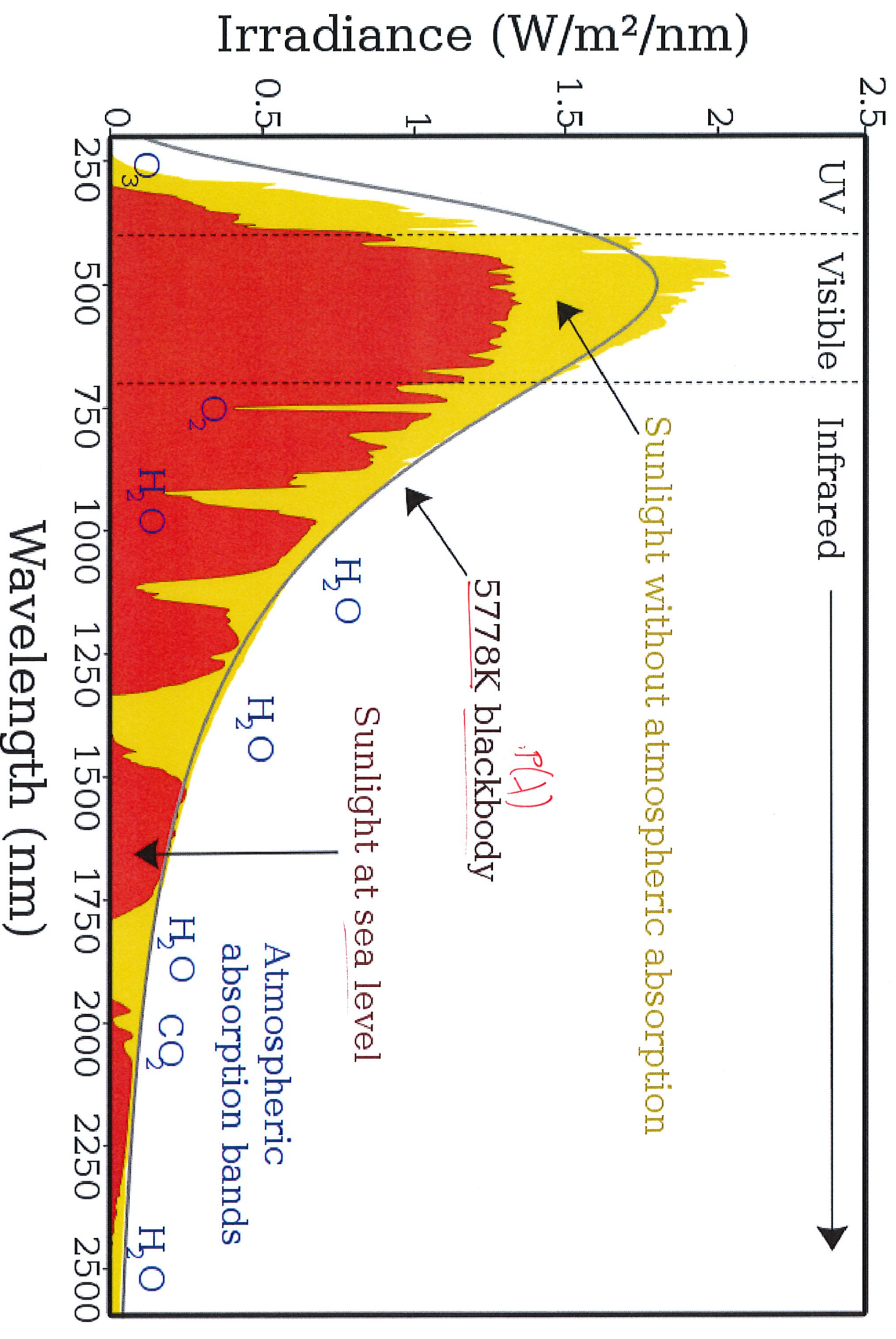
Blue stars $\rightarrow T \approx 10,000$ K

Intergalactic space $\rightarrow T \approx 2.725$ K

Spectral radiance ($\text{kW} \cdot \text{sr}^{-1} \cdot \text{m}^{-2} \cdot \text{nm}^{-1}$)



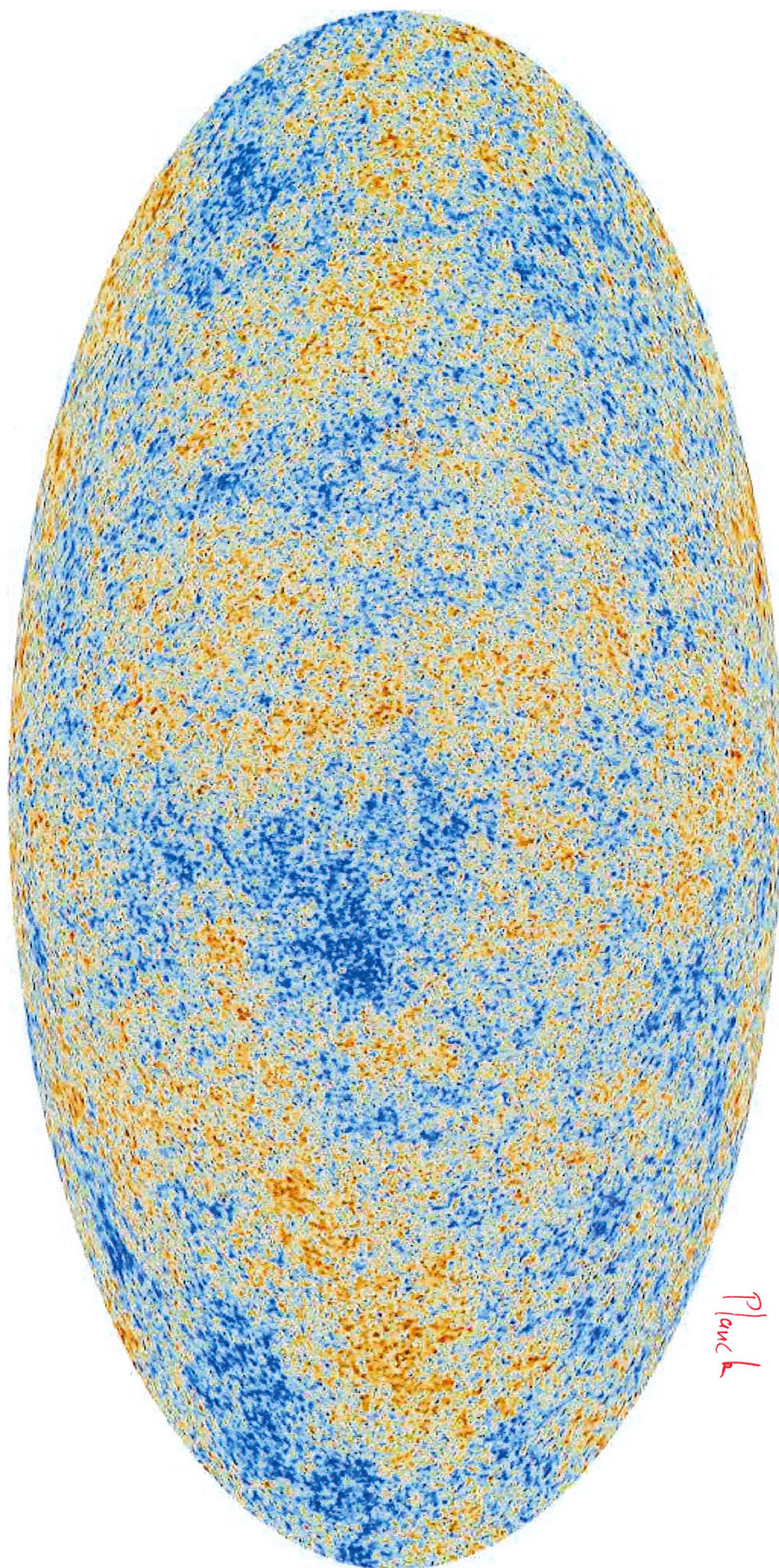
Spectrum of Solar Radiation (Earth)



Cosmic Microwave Background left over from Big Bang ~ 146 yr
T at each point in sky from photons after subtracting galaxies
Blue/red small Fluctuations around average Temp ≈ 2.725 K
 $\Delta T \approx 0.0002$ K

Pattern of fluctuations $\rightarrow$ dark matter

Amazingly accurate descriptions given assumption of non-interacting
ideal gas



Planch

